Bilbao webinar 20/05/2020

CURVATURE FLUCTUATIONS FROM DISORDER DURING INFLATION

Marcos A. G. García

IFT - UAM

1902.09598, 2001.09158
with M. Amin, D. Green, S. Carlsten

Instituto de Física Teórica UAM-CSIC

Universidad Autónoma de Madrid

CSIC Consejo Superior de Investigaciones Científicas

Credit: D. Baumann

Does the simplicity of the data reflect the simplicity of the underlying theory, or does it emerge from complexity?

Complexity in the early universe

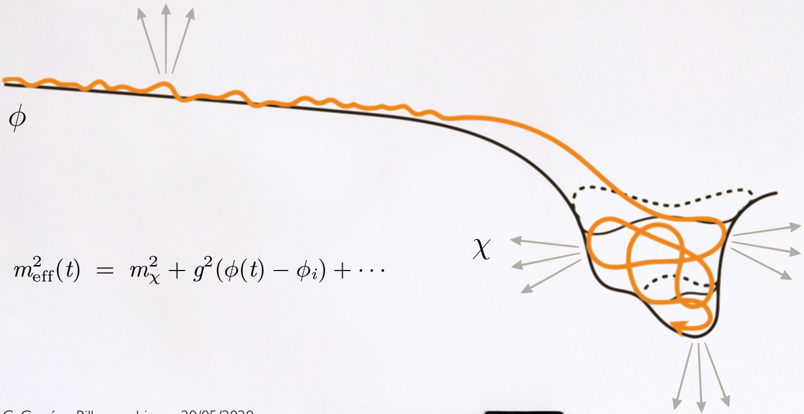
Inflation

- Near scale invariant: $\Delta_\zeta^2 \sim k^{n_s-1}$
- Near Gaussian
- Weak self-interaction (slow roll)

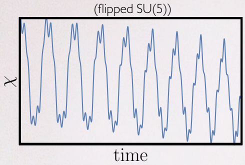
$\longleftrightarrow$

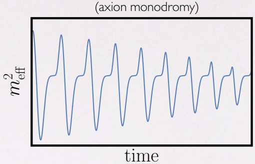
Particle theory

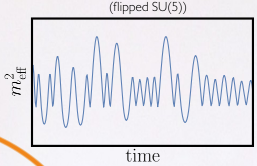
- SM UV completions $N_F \gg 1$
- Coupling to ϕ weakly constrained
- Non-trivial field manifolds

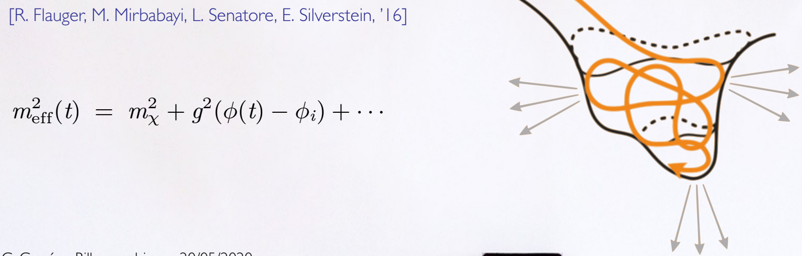


$$m_{\rm eff}^2(t) \;=\; m_\chi^2 + g^2(\phi(t) - \phi_i) + \cdots$$









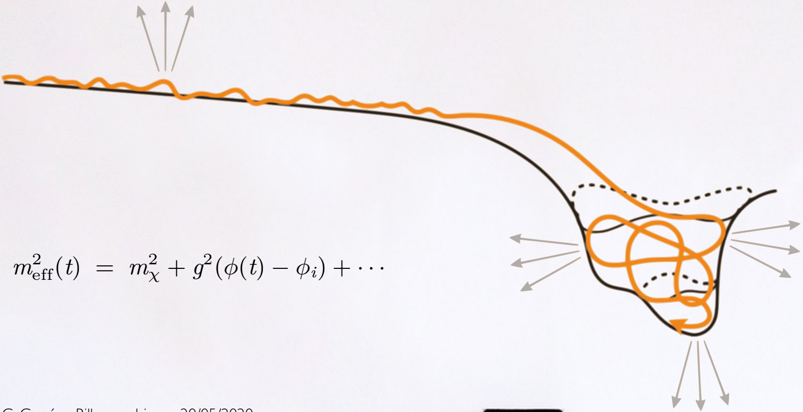
[J. Ellis, MG, N. Nagata, D. Nanopoulos, K. Olive, '17]

[R. Flauger, M. Mirbabayi, L. Senatore, E. Silverstein, '16]

$$m^2_{\rm eff}(t) \;=\; m^2_\chi + g^2(\phi(t) - \phi_i) + \cdots$$

Complexity in the early universe

background dynamics $\longrightarrow$ particle production $\langle \chi_{\boldsymbol{k}_1} \chi_{\boldsymbol{k}_2} \cdots \rangle$ $\longleftrightarrow$ curvature fluctuations $\langle \zeta_{\boldsymbol{k}_1} \zeta_{\boldsymbol{k}_2} \cdots \rangle$



$$m_{\rm eff}^2(t) \;=\; m_\chi^2 + g^2(\phi(t) - \phi_i) + \cdots$$

$$\mathcal{S} = \frac{1}{2}\int \sqrt{-g}\, d^4x \left[c(t+\pi)\partial_\mu\pi\partial^\mu\pi + \partial_\mu\chi\partial^\mu\chi - \left(M^2 + m^2(t+\pi)\right)\chi^2 \right]$$

$$\simeq \frac{1}{2}\int \sqrt{-g}\, d^4x \left[c(t)\partial_\mu\pi\partial^\mu\pi + \partial_\mu\chi\partial^\mu\chi - \left(M^2 + m^2(t)\right)\chi^2 - \frac{dm^2}{dt}\chi^2\pi + \cdots \right]$$

$c = 2M_P^2|\dot{H}|$

Conformal Spectator in dS

$$\mathcal{S} = \frac{1}{2}\int \sqrt{-g}\, d^4x \left[c(t+\pi)\partial_\mu\pi\partial^\mu\pi + \partial_\mu\chi\partial^\mu\chi - \left(M^2 + m^2(t+\pi)\right)\chi^2 \right]$$

$$\simeq \frac{1}{2}\int \sqrt{-g}\, d^4x \left[c(t)\partial_\mu\pi\partial^\mu\pi + \partial_\mu\chi\partial^\mu\chi - \left(M^2 + m^2(t)\right)\chi^2 - \frac{dm^2}{dt}\chi^2\pi + \cdots \right]$$

$$M^2 = 2H^2$$

$$\downarrow$$

$$X_k \equiv a\,\chi_k$$

$$= \alpha_{k,j}\frac{e^{-ik\tau}}{\sqrt{2k}} + \beta_{k,j}\frac{e^{ik\tau}}{\sqrt{2k}}$$

$$m^2(t) = \textstyle\sum_j m_j \delta(t - t_j)$$

$$\langle m_j \rangle = 0$$

$$\langle m_i m_j \rangle = \sigma^2 \delta_{ij}$$

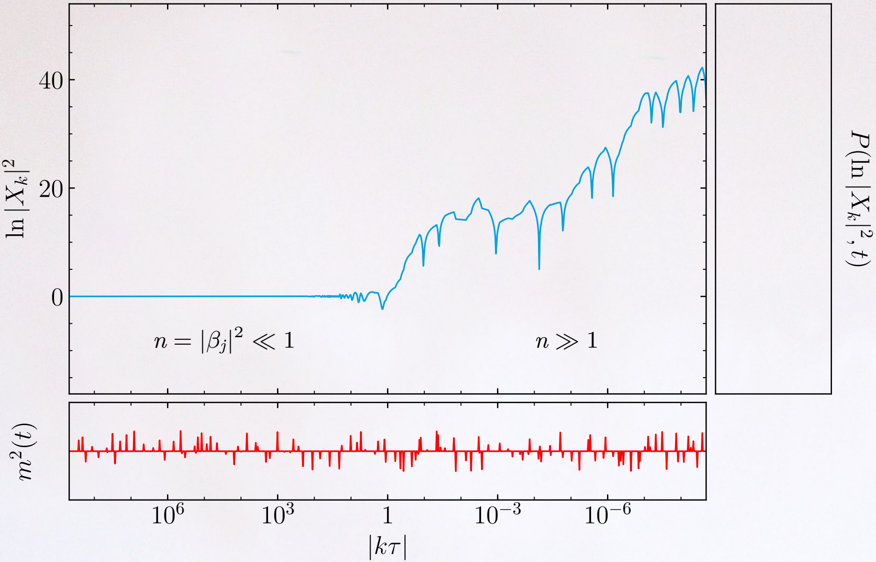
$$\mathcal{N}_s \equiv \frac{\langle N_s \rangle}{N_e}$$

$$\chi \;\underset{\not\leftarrow}{\rightarrow}\; \pi$$

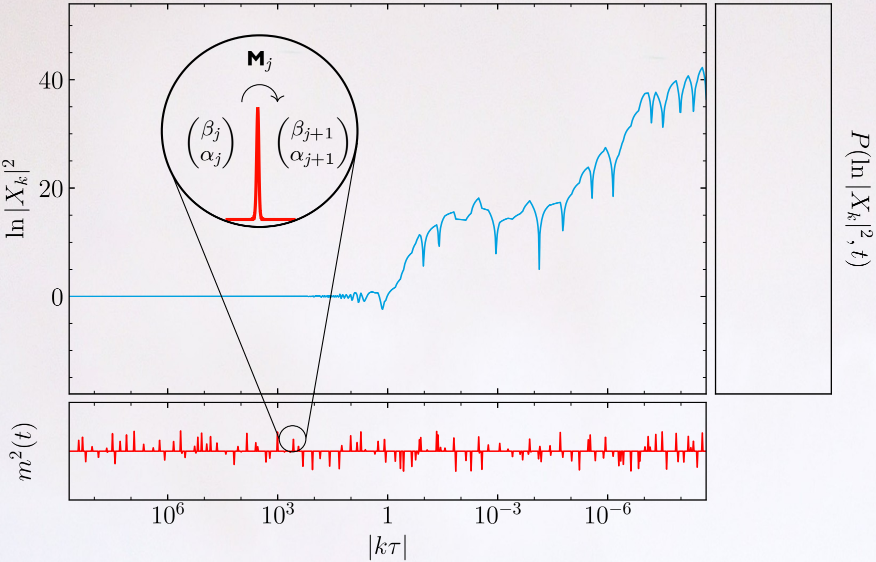
$$\mathcal{O}_S \sim \langle \chi^4 \rangle$$

$$\mathcal{O}_D \sim \langle [\chi^2, \chi^2] \rangle$$

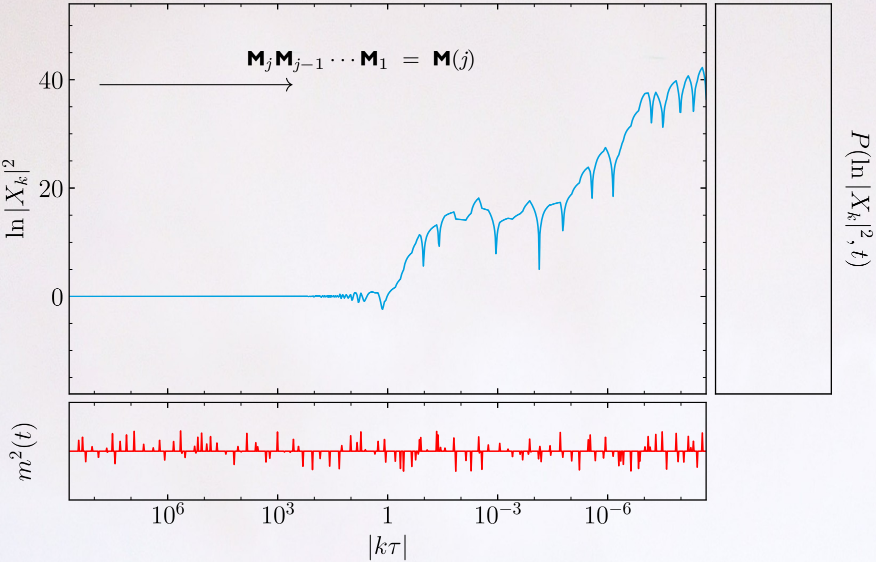
Conformal Spectator in dS

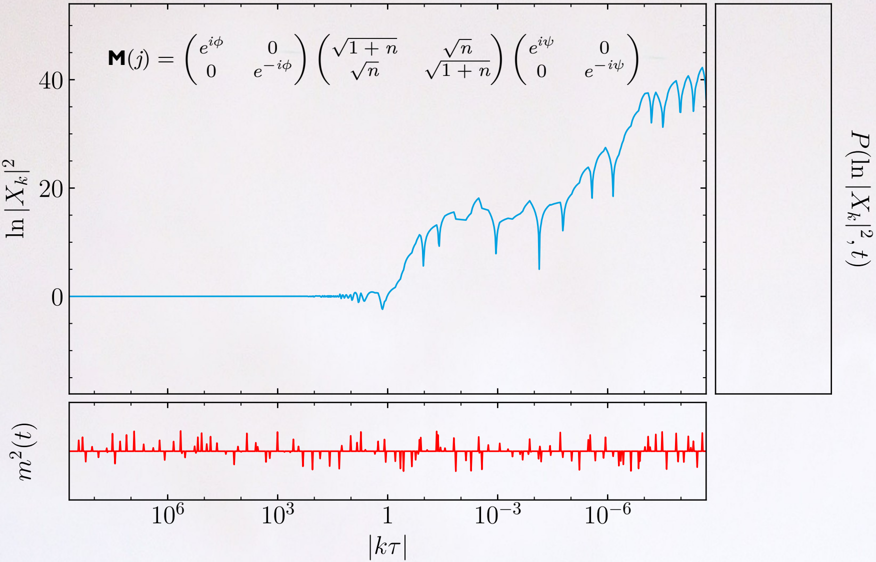


Conformal Spectator in dS

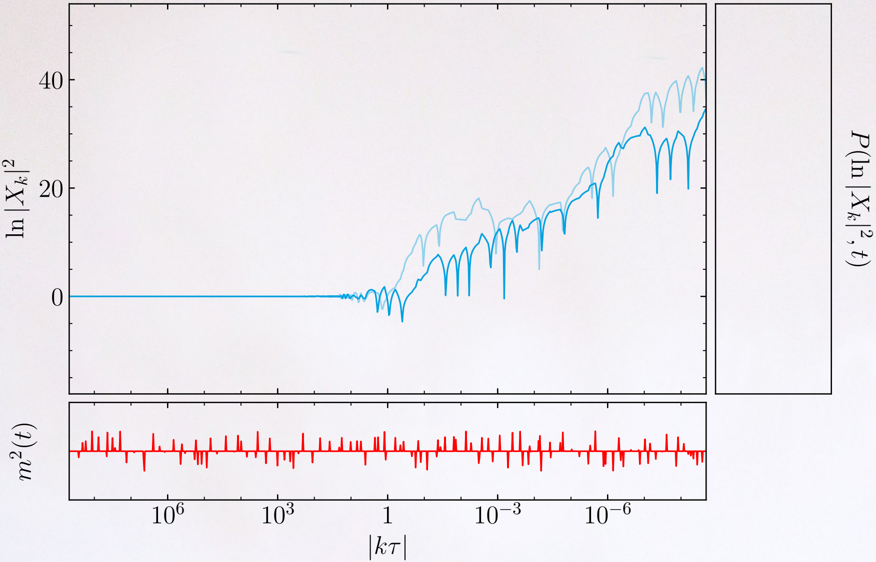


Conformal Spectator in dS

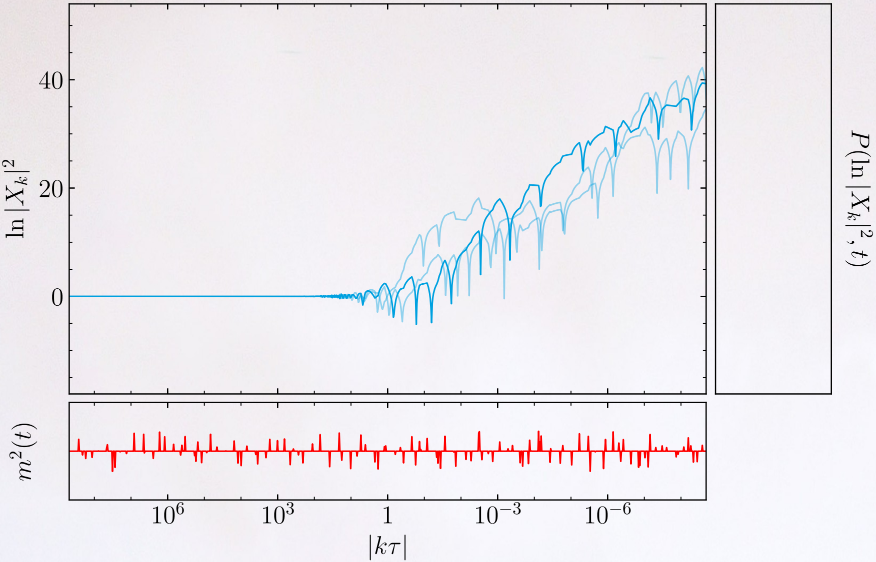


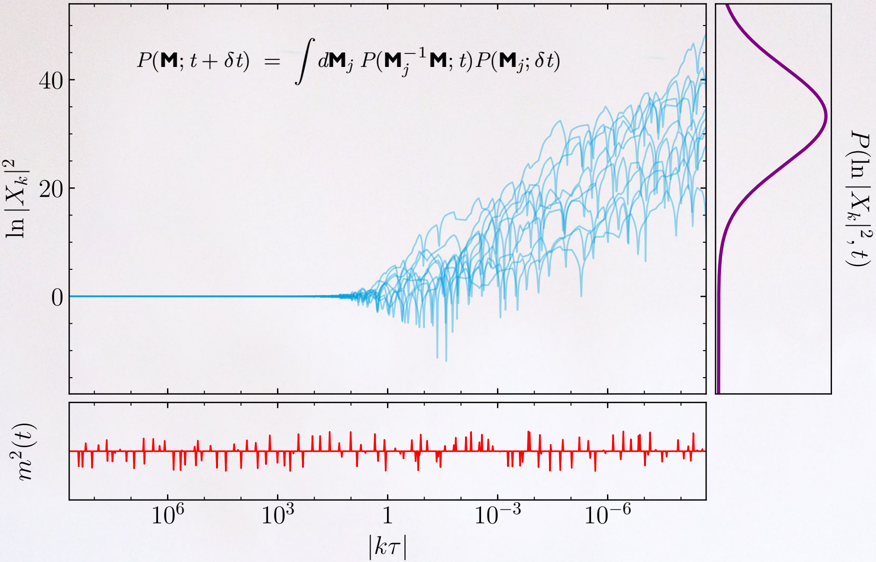

$\mathbf{M}(j) = \begin{pmatrix} e^{i\phi} & 0 \\ 0 & e^{-i\phi} \end{pmatrix} \begin{pmatrix} \sqrt{1+n} & \sqrt{n} \\ \sqrt{n} & \sqrt{1+n} \end{pmatrix} \begin{pmatrix} e^{i\psi} & 0 \\ 0 & e^{-i\psi} \end{pmatrix}$
$\ln |X_k|^2$
40
20
0
$P(\ln |X_k|^2, t)$
$m^2(t)$
10^6
10^3
1
10^{-3}
10^{-6}
$|k\tau|$

Conformal Spectator in dS



Conformal Spectator in dS





$P(\mathbf{M}; t+\delta t) = \int d\mathbf{M}_j\, P(\mathbf{M}_j^{-1}\mathbf{M}; t) P(\mathbf{M}_j; \delta t)$
40
20
0
$\ln|X_k|^2$
$m^2(t)$
10^6
10^3
1
10^{-3}
10^{-6}
$|k\tau|$
$P(\ln|X_k|^2, t)$

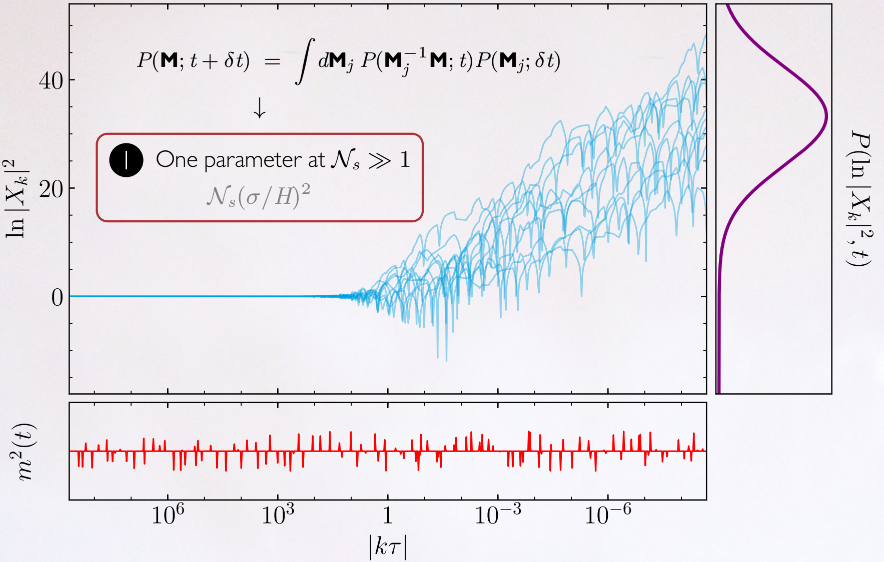
Conformal Spectator in dS

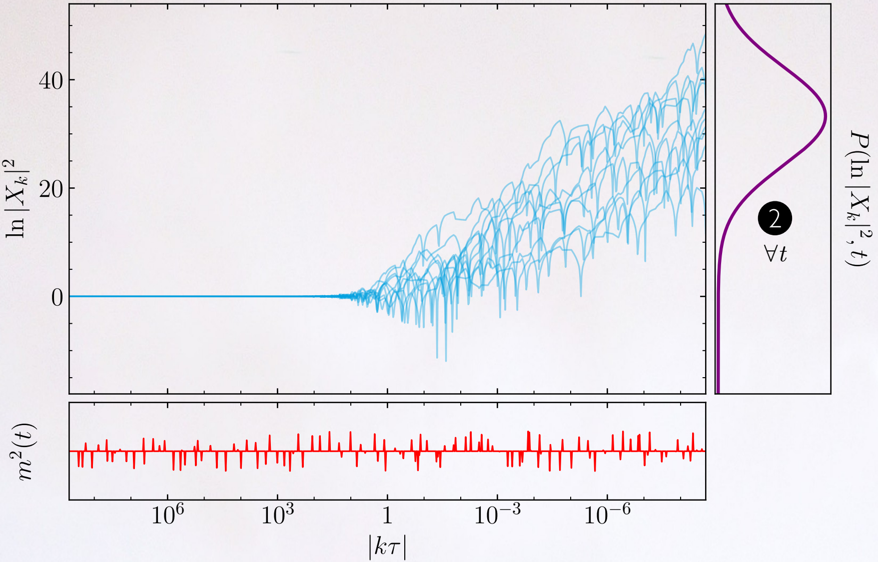
$$P(\mathbf{M}; t+\delta t) = \int d\mathbf{M}_j \, P(\mathbf{M}_j^{-1}\mathbf{M}; t) P(\mathbf{M}_j; \delta t)$$

$\downarrow$

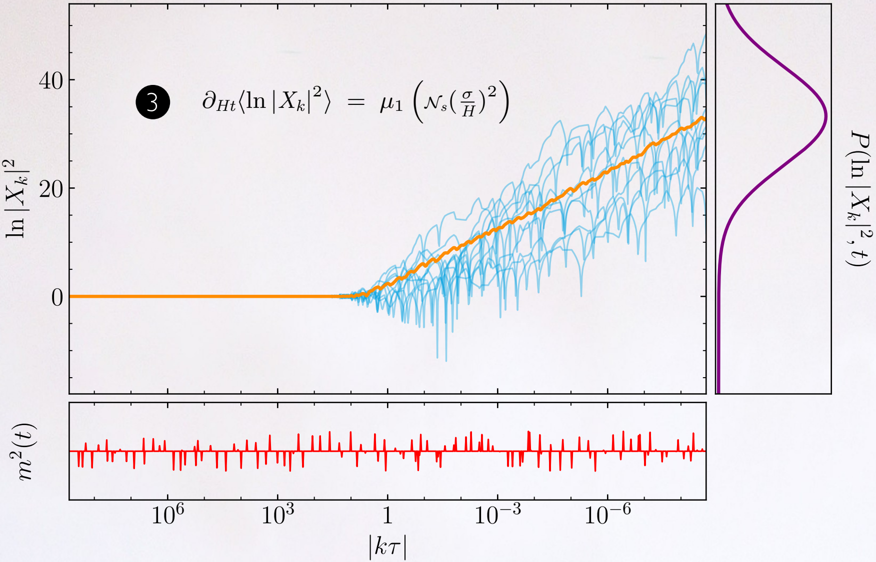
I One parameter at $\mathcal{N}_s \gg 1$

$\mathcal{N}_s(\sigma/H)^2$



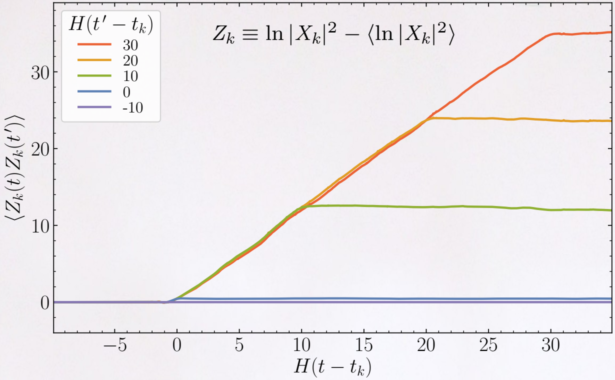


$\ln|X_k|^2$
40
20
0
$m^2(t)$
10^6
10^3
1
10^{-3}
10^{-6}
$|k\tau|$
$P(\ln|X_k|^2, t)$
2
$\forall t$



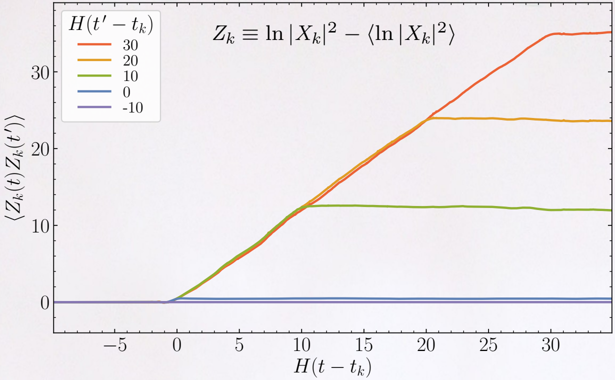
3
$\partial_{Ht}\langle \ln |X_k|^2\rangle \;=\; \mu_1 \left(\mathcal{N}_s\left(\frac{\sigma}{H}\right)^2 \right)$
40
20
0
$\ln |X_k|^2$
$P(\ln |X_k|^2, t)$
$m^2(t)$
10^6
10^3
1
10^{-3}
10^{-6}
$|k\tau|$

4



$$\langle Z_k(t) Z_{k'}(t') \rangle \;=\; \mu_2 \left(\mathcal{N}_s \left(\frac{\sigma}{H}\right)^2 \right) H \min\left[t - t_k, t - t_{k'}, t' - t_k, t' - t_{k'}\right]$$

❶ + ❷ + ❸ + ❹ = geometric (Brownian) random walk of $|X_k|^2$



$$\left\langle |X_{k_1}(t_1)|^2 \cdots |X_{k_n}(t_n)|^2 \right\rangle = \exp\left[\sum_{i=1}^{n} \langle \ln |X_{k_i}(t_i)|^2 \rangle + \frac{1}{2} \sum_{i,j=1}^{n} \langle Z_{k_i}(t_i) Z_{k_j}(t_j) \rangle \right]$$

Curvature power spectrum

To lowest order in π, with $\zeta \simeq H\pi$ and $\langle\zeta(\mathbf{k})\zeta(\mathbf{k}')\rangle = \frac{2\pi^2}{k^3}\Delta_\zeta^2(k)\,\delta^{(3)}(\mathbf{k}+\mathbf{k}')$,

$$\hat{\pi}''(\mathbf{x},\tau) + 2\mathcal{H}\hat{\pi}'(\mathbf{x},\tau) - \nabla^2\hat{\pi}(\mathbf{x},\tau) \;=\; -\frac{a(\tau)}{2c(\tau)}\frac{dm^2(\tau)}{d\tau}\hat{\chi}^2(\mathbf{x},\tau)\,,$$

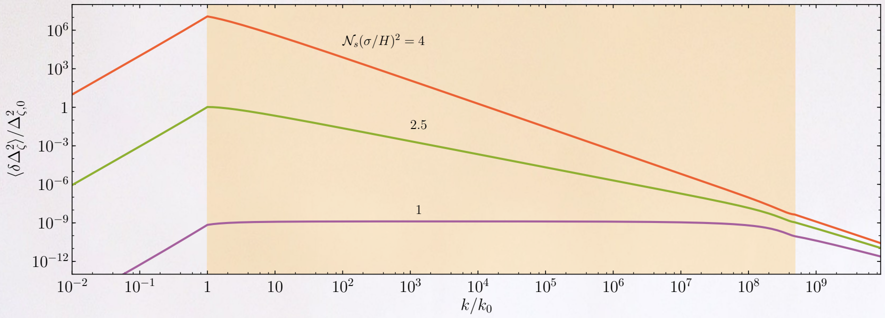
$$\downarrow$$

$$\delta\Delta_\zeta^2(k) \;=\; 4\pi^2(\Delta_{\zeta,0}^2)^2\frac{k^3}{H^4}\int d\tau'\,d\tau''\,\tau'\tau''\,G_k(\tau,\tau')\,G_k(\tau,\tau'')\,\frac{dm^2(\tau')}{d\tau'}\,\frac{dm^2(\tau'')}{d\tau''}$$
$$\times\int^{\Lambda(\tau)}\frac{d^3\mathbf{p}}{(2\pi)^3}\,\left[X_p(\tau')X_p^*(\tau'')\right]_{\rm AS}\left[X_{|\mathbf{p}-\mathbf{k}|}(\tau')X_{|\mathbf{p}-\mathbf{k}|}^*(\tau'')\right]_{\rm AS}$$

- $\mathcal{O}_{\rm AS} \;\equiv\; \mathcal{O} - \mathcal{O}_{\rm vac}$
- $\Lambda(\tau) \;\equiv\; (Hw\tau)^{-1}$
- Start: $|k_0\tau| = 1$. End: $|k_f\tau| = 1$

Curvature power spectrum

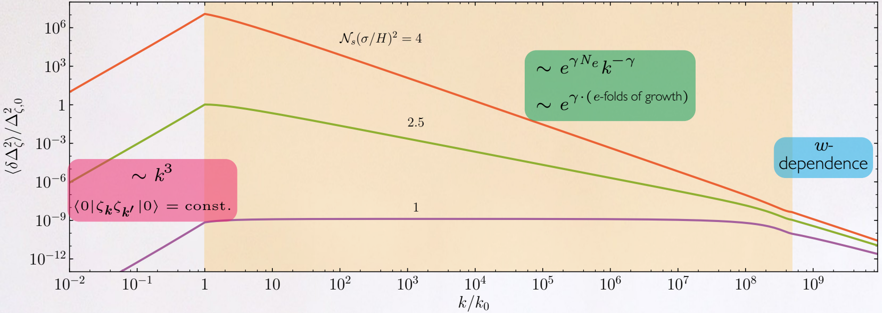
$$\left\langle \delta\Delta_\zeta^2(k) \right\rangle = \left(\Delta_{\zeta,0}^2\right)^2 \mathcal{N}_s \left(\frac{\sigma}{H}\right)^2 e^{\mathcal{F}\left(k, N_e, \mathcal{N}_s(\sigma/H)^2\right)}$$



always super-horizon | cross horizon during scatterings | always sub-horizon

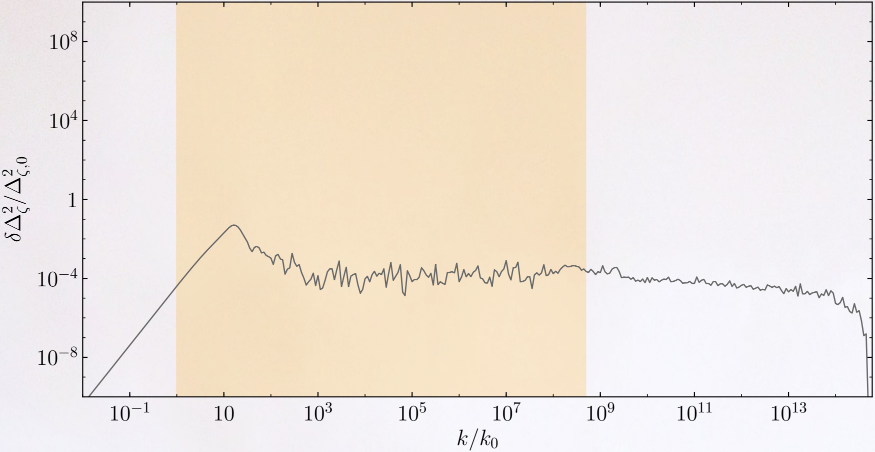
Curvature power spectrum

$$\left\langle \delta\Delta^2_\zeta(k) \right\rangle = \left(\Delta^2_{\zeta,0}\right)^2 \mathcal{N}_s \left(\frac{\sigma}{H}\right)^2 e^{\mathcal{F}\left(k, N_e, \mathcal{N}_s(\sigma/H)^2\right)}$$

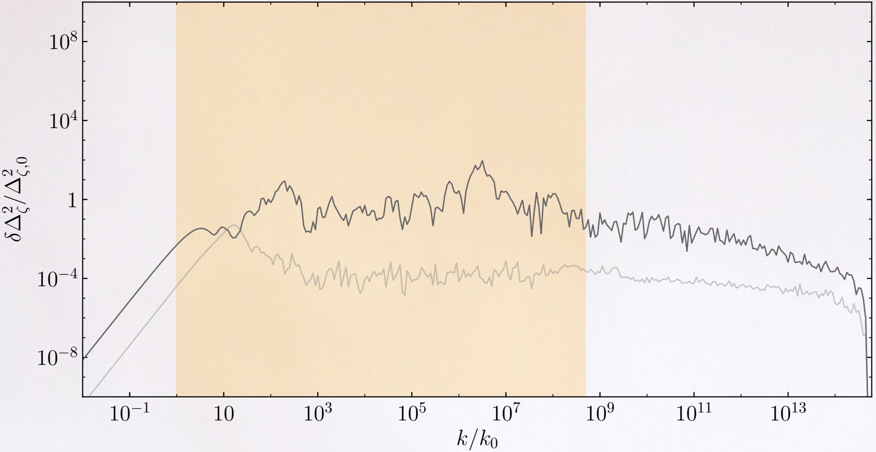


always super-horizon | cross horizon during scatterings | always sub-horizon

$(N_e = 20, \mathcal{N}_s(\sigma/H)^2 = 25)$

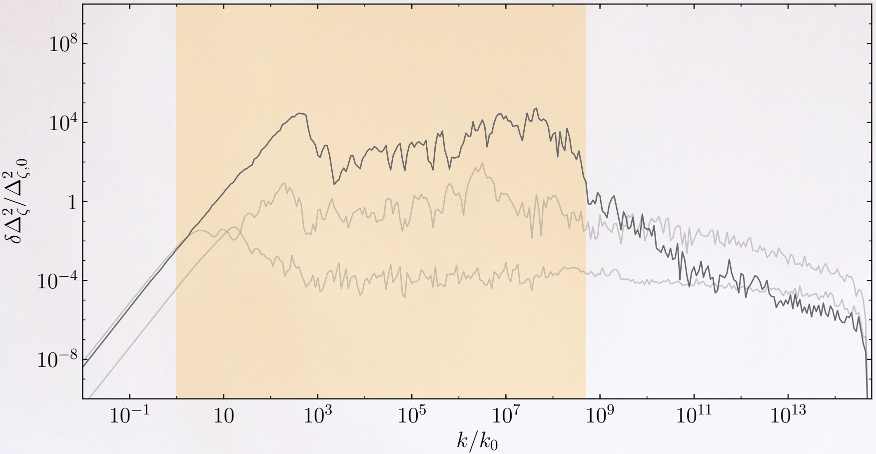


$(N_e = 20, \mathcal{N}_s(\sigma/H)^2 = 25)$



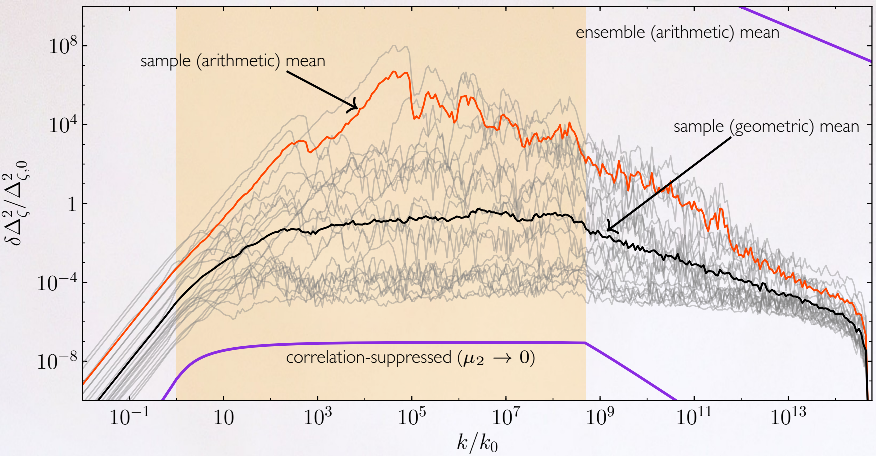
Curvature power spectrum

$(N_e = 20, \mathcal{N}_s(\sigma/H)^2 = 25)$

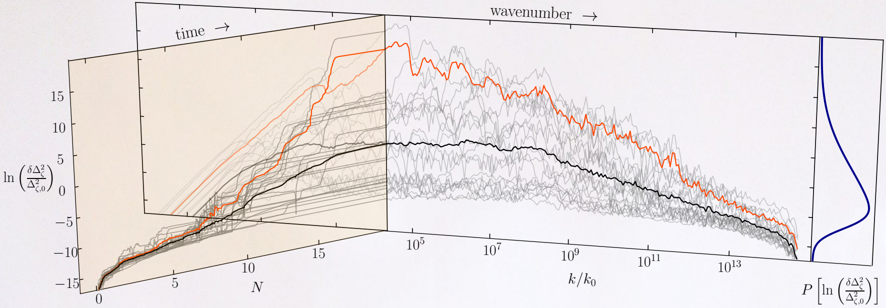


Curvature power spectrum

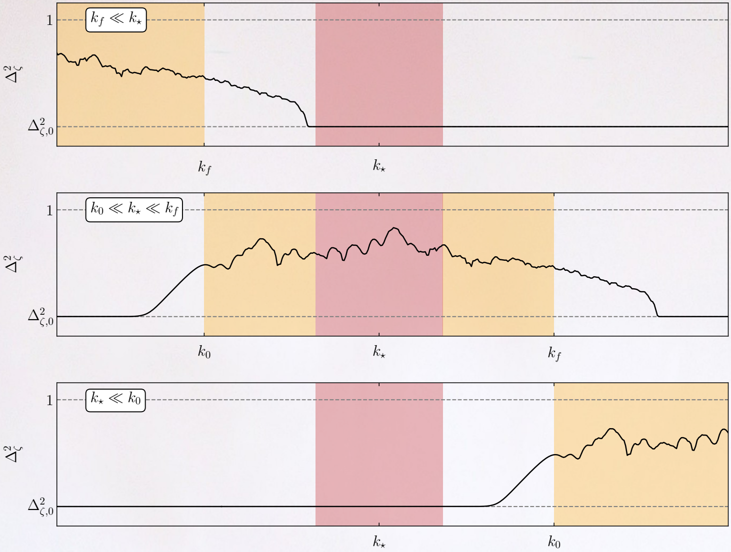
$(N_e = 20, \mathcal{N}_s(\sigma/H)^2 = 25)$



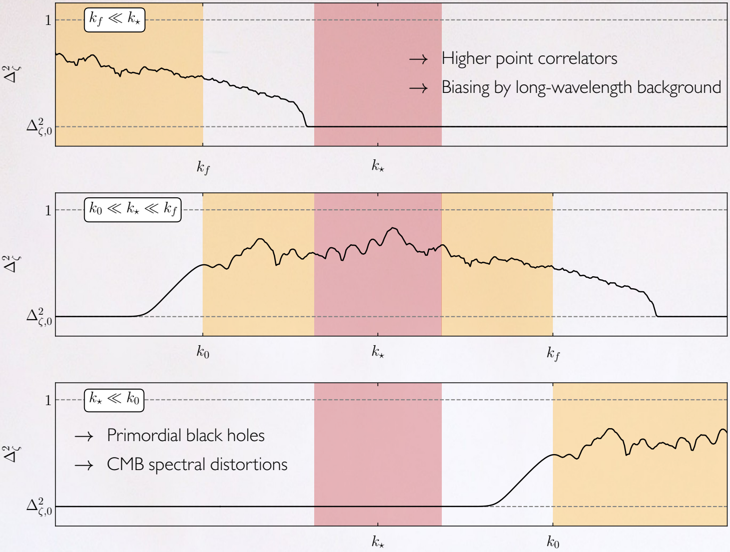
$(N_e = 20, \mathcal{N}_s(\sigma/H)^2 = 25)$



Observational implications

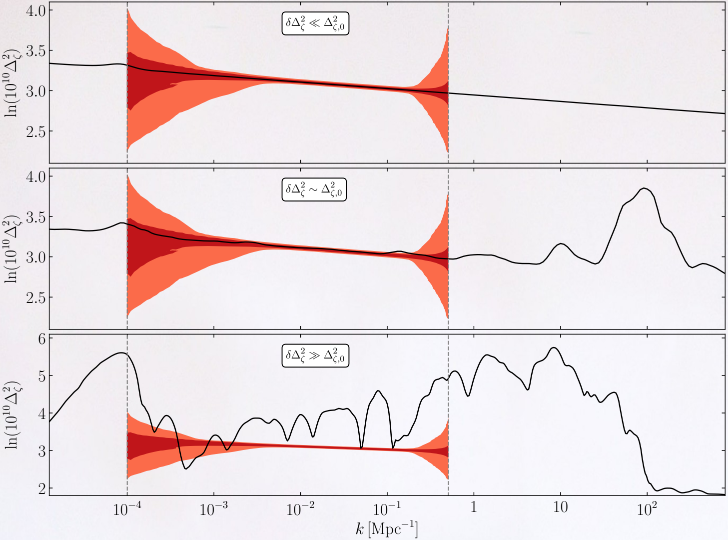


Observational implications



Observational implications

$\mathcal{N}_s(\sigma/H)^2 = 25, \quad N_{\rm tot} = 20,$ Planck TT, TE, EE + lowE + lensing + BK15 (1σ, 2σ)



Conclusions

- Stochastically excited spectator fields undergo geometric random walks
- Lead to features in the curvature power spectrum $\to$ constraints
- Look for enhancement in the N-point function

$$\langle\zeta^n\rangle - \langle\zeta^n\rangle_{\chi=0} \sim \langle\zeta^2\rangle^n_{\chi=0} \times \exp\left[\frac{n^2}{2} F\left(\mathcal{N}_s \frac{\sigma^2}{H^2}\right)\right]$$

- Higher spin spectators / higher spin observables
- Stochastic preheating
- Backreaction regime $\to$ dissipation

Thank You