



Freeze-in from Preheating



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Consider FIMPs:

10-2

Matter-nu Matter-nu

1000

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cleon $\sigma_{\rm SI}$

Dark Dark

- Never in thermal equilibrium
- Produced via freeze-in
- Elusive (in)direct detection
- Dependence on initial conditions (inflation, reheating)



Freeze-in

$$\frac{\partial f_{\rm DM}}{\partial t} - Hp \frac{\partial f_{\rm DM}}{\partial p} = \frac{1}{16p} \int \frac{d^3 \mathbf{p}'}{(2\pi)^3 p'_0} \frac{d^3 \mathbf{k}}{(2\pi)^3 k_0} \frac{d^3 \mathbf{k}'}{(2\pi)^3 k'_0} (2\pi)^4 \delta(\mathbf{p} + \mathbf{p}' - \mathbf{k} - \mathbf{k}') |\mathcal{M}|^2_{\to} f_i(k) f_i(k')$$

$$\frac{dn_{\rm DM}}{dt} + 3Hn_{\rm DM} = \int \frac{d^3 \mathbf{k}}{(2\pi)^3 k_0} \frac{d^3 \mathbf{k}'}{(2\pi)^3 k'_0} (\mathbf{k} \cdot \mathbf{k}') \,\sigma(s) \,f_i(k) f_i(k')$$







4. Relic abundances



 $k^{2}f_{i}(k)$:





4. Relic abundances

The inflaton and its decay products

 \mathcal{S}

$$\begin{split} &= \int d^4x \sqrt{-g} \left[\frac{1}{2} (\partial_\mu \phi)^2 - 6\lambda M_P^4 \tanh^2 \left(\frac{\phi}{\sqrt{6}M_P} \right) & \text{ inflaton} \\ &+ \frac{1}{2} (\partial_\mu \chi)^2 - \frac{1}{2} \sigma \phi^2 \chi^2 & \text{ scalar } (m_\chi^2 = \sigma \phi^2) \\ &+ \bar{\psi} i \bar{\gamma}^\mu \nabla_\mu \psi - y \phi \bar{\psi} \psi + \cdots \right] & \text{ fermion } (m_\psi^2 = y^2 \phi^2) \end{split}$$



$$p_{\phi}\rangle = \frac{1}{2} \langle \dot{\phi}^2 + m_{\phi}^2 \phi^2 \rangle \simeq 0$$
 (matter)



3. Preheating

4. Relic

abundances

The fluid picture

$$T^{\mu\nu} =
ho_{\phi} \operatorname{diag}(1,0,0,0) + rac{
ho_R}{3} \operatorname{diag}(3,1,1,1)$$

$$\nabla_{\mu} T^{\mu\nu} = 0 \qquad \Rightarrow \qquad \dot{\rho}_{R} + 4H\rho_{R} = -(\dot{\rho}_{\phi} + 3H\rho_{\phi}) \equiv \Gamma_{\phi}\rho_{\phi}$$

Decay rate of an oscillating inflaton *condensate*:

 $\phi(t) \simeq \phi_0(t) \mathcal{P}(t) = \phi_0(t) \sum_n \mathcal{P}_n e^{-in\omega t}$

$$\Gamma_{\phi} = \frac{1}{8\pi\rho_{\phi}} \sum_{n=1}^{\infty} E_n \left\langle |\mathcal{M}_n|^2 \beta_n \right\rangle$$

 $E_n = n m_{\phi}$ $\beta_n = \sqrt{1 - \frac{4m_{\chi,\psi}^2}{E_n^2}}$ $\langle f | i \int d^4 x \mathcal{L}_I | 0 \rangle = i(2\pi)^4 \sum_{n=-\infty}^{\infty} \mathcal{M}_n \, \delta^{(4)}(p_n - p_1 - p_2)$

(For $V \propto \phi^k$ see MG, K. Kaneta, Y. Mambrini, K. Olive, JCAP 04 (2021), 012)



Kinematic blocking: scalars

 $\Gamma_{\phi\phi\to\chi\chi} = \frac{\sigma^2 \rho_{\phi}(t)}{32\pi m_{\star}^3} \alpha_{\sigma}(\mathcal{R})$





3. Preheating



4. Relic abundances







2. Perturbative reheating



3. Preheating



4. Relic abundances

Kinematic blocking: fermions

$$\beta_n = \left(1 - \frac{\mathcal{R}}{n^2} \mathcal{P}^2\right)^{1/2}$$

$${\cal R} \;=\; 2 {y^2 \over \lambda} \left({\phi_0 \over M_P}
ight)^2$$











Scalar preheating

$$\ddot{\chi}_p + 3H\dot{\chi}_p + \left[\frac{p^2}{a^2} + m_{\chi}^2(t)\right]\chi_p = 0, \qquad m_{\chi}^2(t) = \frac{1}{2}\sigma\phi^2$$

With expansion,





2. Perturbative

reheating

3. Preheating

Unstable-scalar preheating

 $\left(\frac{d^2}{dt^2} - \frac{\nabla^2}{a^2} + (3H + \Gamma_\chi)\frac{d}{dt} + m_\chi^2(t)\right)\chi = 0 \qquad \Gamma_\chi = \frac{y_\chi^2}{8\pi}m_\chi(t)$ $\dot{\rho}_f + 4H\rho_f = \Gamma_\chi\dot{\chi}^2$

Or, with
$$Y_p \equiv a \exp\left(\frac{1}{2} \int a\Gamma_{\chi} d\tau\right) \chi_p$$

$$Y''_p + \omega_p^2 Y_p = 0$$
$$\rho'_f + 4\mathcal{H}\rho_f = a\Gamma_{\chi}\dot{\chi}^2$$

where

4. Relic abundances

$$\omega_p^2 \equiv p^2 + a^2 m_\chi^2 - rac{a^{\prime\prime}}{a} - rac{1}{4} (a \Gamma_\chi)^2 - rac{3}{2} a \mathcal{H} \Gamma_\chi$$

$$\rho_{\chi} = \frac{e^{-\int a\Gamma_{\chi} d\tau}}{(2\pi)^3 a^4} \int d^3 \boldsymbol{p} \, \omega_p n_p \,, \qquad n_p = \frac{1}{2\omega_p} |\omega_p \, Y_p - i Y_p'|^2$$











Fermion preheating

$$\left[i\gamma^{\mu}\partial_{\mu} + i\frac{3a'}{2a}\gamma^{0} - am_{\psi}(\tau)\right]\psi = 0, \qquad m_{\psi}^{2}(\tau) = y^{2}\phi^{2}$$

2. Perturbative reheating





4. Relic abundances

Neglecting expansion,*

$$^{st}dt/d au=a$$
 and

with

$$\psi(\tau, \mathbf{x}) = a^{-3/2} \sum_{r=\pm} \int \frac{d^3 \mathbf{p}}{(2\pi)^{3/2}} e^{-i\mathbf{p}\cdot\mathbf{x}} \left[u_p^{(r)}(\tau) \hat{a}_{\mathbf{p}}^{(r)} + v_p^{(r)}(\tau) \hat{b}_{-\mathbf{p}}^{(r)\dagger} \right]$$

$$u_p^{(r)}(au) \;=\; egin{pmatrix} U_1(au) \xi_r(oldsymbol{p}) \ U_2(au) rac{oldsymbol{\sigma} \cdot oldsymbol{p}}{p} \xi_r(oldsymbol{p}) \end{pmatrix}$$



Fermion preheating

$$\left[i\gamma^{\mu}\partial_{\mu} + i\frac{3a'}{2a}\gamma^{0} - am_{\psi}(\tau)\right]\psi = 0, \qquad m_{\psi}^{2}(\tau) = y^{2}\phi^{2}$$

With expansion,





10

.

 10^{2}



Freeze-in from scalar preheating

Reminder:

with

2. Perturbative reheating



3. Preheating



4. Relic abundances

$$\frac{dn_{\rm DM}}{dt} + 3Hn_{\rm DM} = \int \frac{d^3 \mathbf{k}}{(2\pi)^3 k_0} \frac{d^3 \mathbf{k}'}{(2\pi)^3 k_0'} \sqrt{(k \cdot k')^2 - m_\chi^2} \,\sigma(s) \, f_\chi(k) f_\chi(k')$$

$$|\mathcal{M}|^2 = 16\pi \frac{s^{\frac{n}{2}+1}}{\Lambda^{n+2}}$$

Parent distributions

Perturbative:

 $\frac{\partial f_{\chi}}{\partial t} - HP \frac{\partial f_{\chi}}{\partial P} = \frac{\pi^2}{\beta^2 m_{\phi}^3} \rho_{\phi} \Gamma_{\phi\phi\to\chi\chi} (1 + 2f_{\chi}) \delta \left(P - \beta m_{\phi}\right)$

Non-perturbative: $f_{\chi}(P,t) = \frac{1}{2\omega_p} \left| \omega_p X_p - i X'_p \right|^2_{p \to aP}$









unstable

 10^{3}

 σ/λ

 $y_{\chi}=1\,,~\Lambda=\overline{10^{14}\,{\rm GeV}}$

 10^{4}

 10^{5}



Freeze-in from fermion preheating

Reminder:

2. Perturbative reheating



3. Preheating



4. Relic abundances

$$\frac{dn_{\rm DM}}{dt} + 3Hn_{\rm DM} = \int \frac{d^3 \mathbf{k}}{(2\pi)^3 k_0} \frac{d^3 \mathbf{k}'}{(2\pi)^3 k_0'} \sqrt{(k \cdot k')^2 - m_{\psi}^2} \,\sigma(s) \, f_{\psi}(k) f_{\psi}(k')$$

 $|\mathcal{M}|^2 = 16\pi \frac{s^{\frac{n}{2}+1}}{\Lambda^{n+2}}$

with

Parent distributions

Perturbative:

Perturbative:
$$\frac{\partial f_{\psi}}{\partial t} - HP \frac{\partial f_{\psi}}{\partial P} = \frac{8\pi^2}{\beta^2 m_{\phi}^3} \rho_{\phi} \Gamma_{\phi \to \bar{\psi}\psi} (1 - 2f_{\psi}) \delta \left(P - \frac{1}{2}\beta m_{\phi}^2\right)^2$$
Non-perturbative:
$$f_{\psi}(P, t) = \frac{1}{2} \left| \left(1 + \frac{m_{\psi}}{\omega_p}\right)^{1/2} U_2 - \left(1 - \frac{m_{\psi}}{\omega_p}\right)^{1/2} U_1 \right|_{p \to aP}^2$$



Fermion PSD











abundances

Backreaction

Homogeneous limit: Hartree approximation

 $\ddot{\phi} + 3H\dot{\phi} + V'(\phi) + \sigma \langle \chi^2 \rangle \phi = 0$

Fragmentation: classical fields in the lattice

D. Figueroa, et al., arXiv:2102.01031 [astro-ph.CO]

 $\Box \phi + \Gamma_{\phi} u^{\mu} \nabla_{\mu} \phi + V_{\phi}(\phi, \chi) = 0$







Lyman- α constraints









Fermion rates



3. Preheating



4. Relic abundances

